

# Chapter 29: Electromagnetic Waves

## Thursday November 10<sup>th</sup>

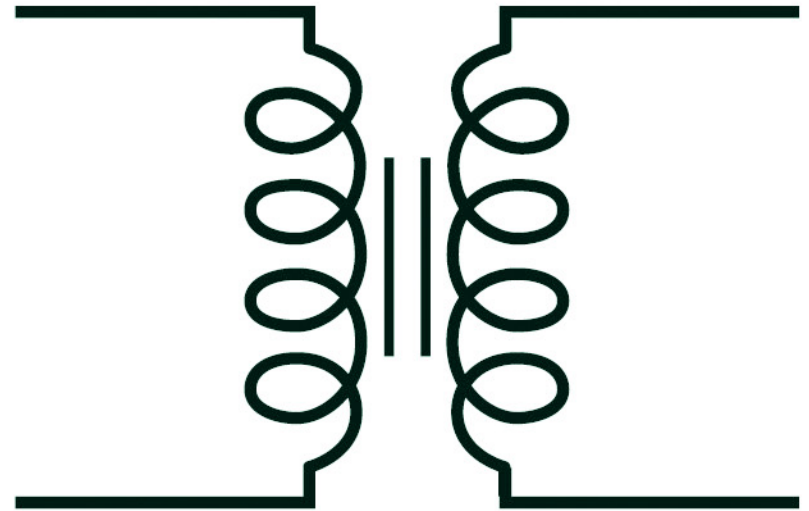
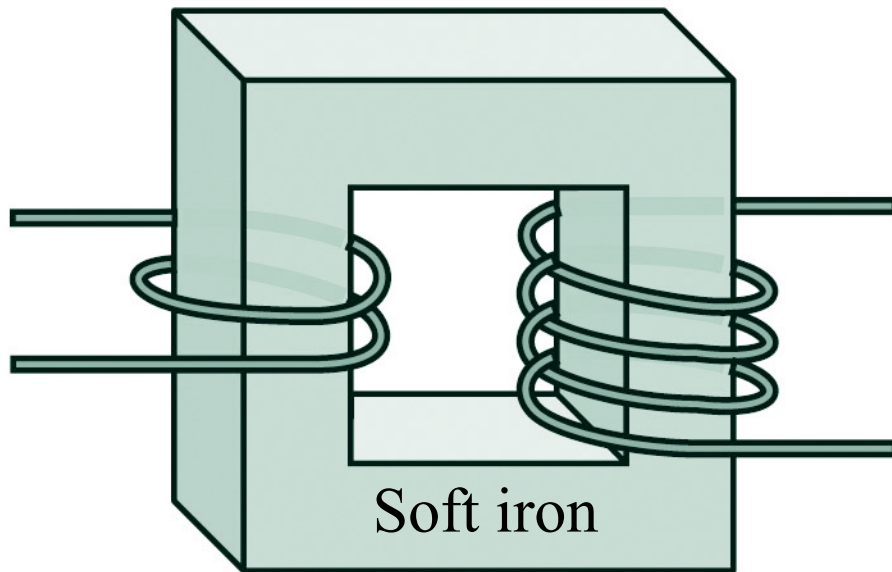
**Mini-exam 5 next Thursday (AC circuits and EM waves)**  
**55 unregistered *iClickers* – any takers?**

- Transformers - demo
- Maxwell's equations
- Electromagnetic waves
  - Wave equations
  - Speed of light
  - Relations between quantities
  - Energy flux and intensity

**Reading: up to page 515 in the text book (Ch. 28/29)**

# Transformers

Primary      Secondary



- Flux the same on both sides, but number of turns,  $N$ , is different
- Total flux through primary and secondary coils depends on  $N_1$  and  $N_2$

$$V_1 = N_1 \Phi; \quad V_2 = N_2 \Phi; \quad \Rightarrow \frac{V_2}{V_1} = \frac{N_2}{N_1}$$

# The basic equations of electromagnetism so far.....

**Gauss' law:**

$$\Phi_E = \oint \vec{E} \cdot d\vec{A} = \frac{q_{enc}}{\epsilon_0}$$

**Gauss' law for B (no magnetic equivalent of charge):**

$$\Phi_B = \oint \vec{B} \cdot d\vec{A} = 0$$

**Ampère's law:**

$$\oint \vec{B} \cdot d\vec{l} = \mu_0 I_{enc}$$

**Faraday's law:**

$$\epsilon = \oint \vec{E} \cdot d\vec{l} = -\frac{d\Phi_B}{dt}$$

# The basic equations of electromagnetism so far.....

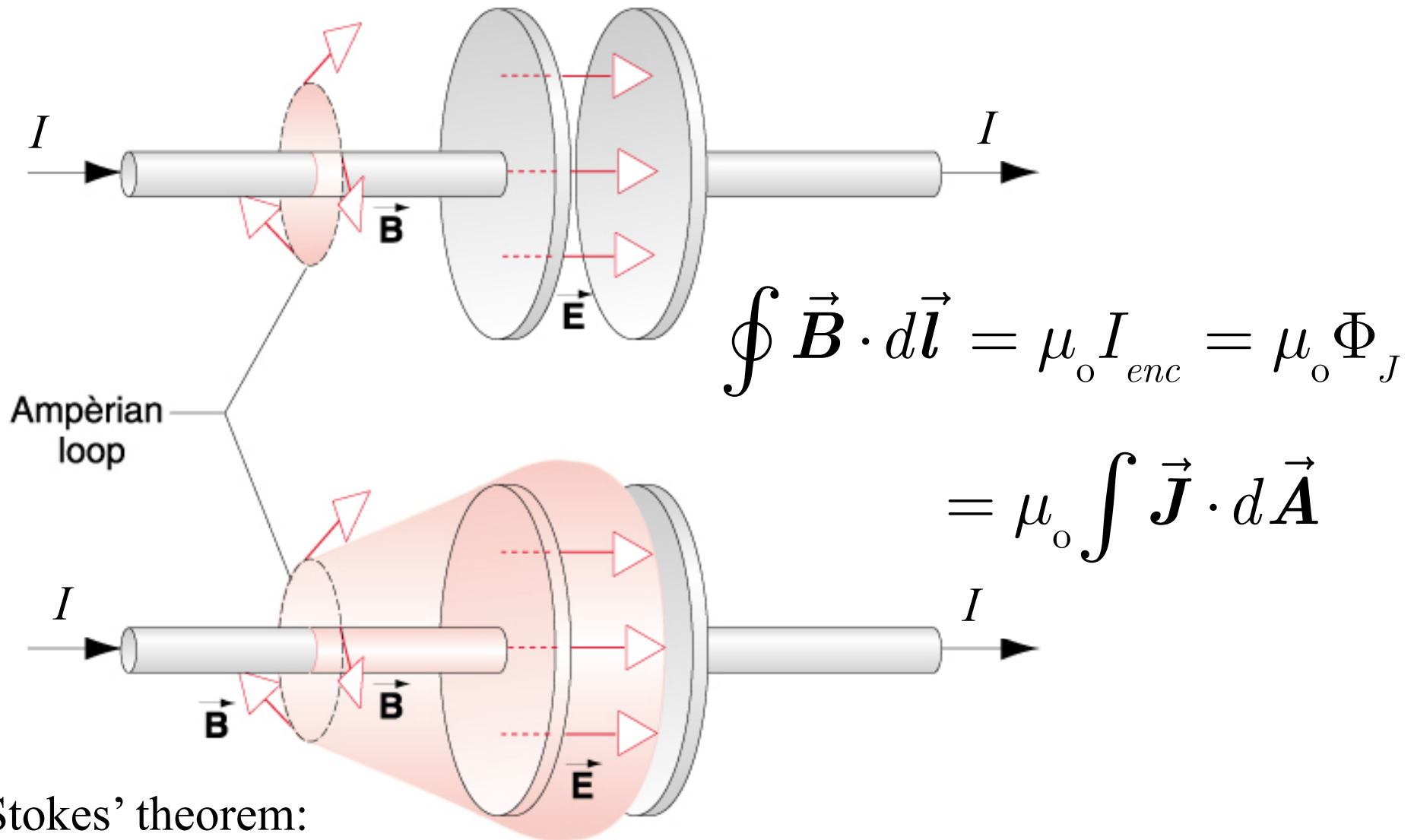
In vacuum:

$$\begin{aligned}\Phi_E &= \oint \vec{E} \cdot d\vec{A} = 0 \\ \Phi_B &= \oint \vec{B} \cdot d\vec{A} = 0\end{aligned} \quad \left. \vphantom{\begin{aligned}\Phi_E \\ \Phi_B\end{aligned}} \right\} \text{Symmetry}$$

$$\begin{aligned}\oint \vec{B} \cdot d\vec{l} &= 0 \\ \oint \vec{E} \cdot d\vec{l} &= -\frac{d\Phi_B}{dt}\end{aligned} \quad \left. \vphantom{\begin{aligned}\oint \vec{B} \cdot d\vec{l} \\ \oint \vec{E} \cdot d\vec{l}\end{aligned}} \right\} \text{No Symmetry}$$

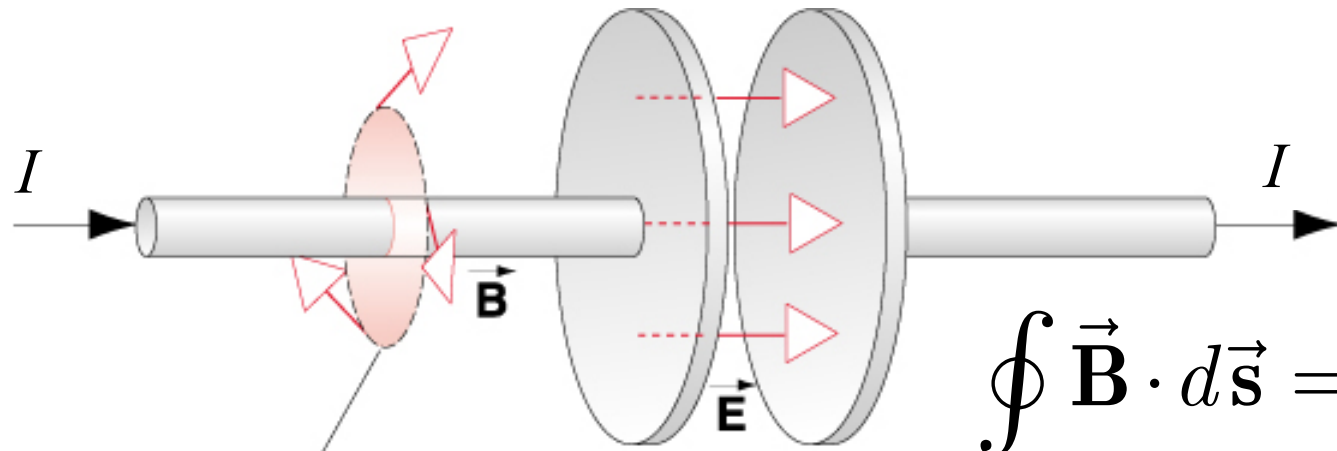
Is it possible that a time-varying electric field could produce a magnetic field, thereby restoring symmetry?

# Maxwell's displacement current



Stokes' theorem:  
The choice of surface  
should not matter

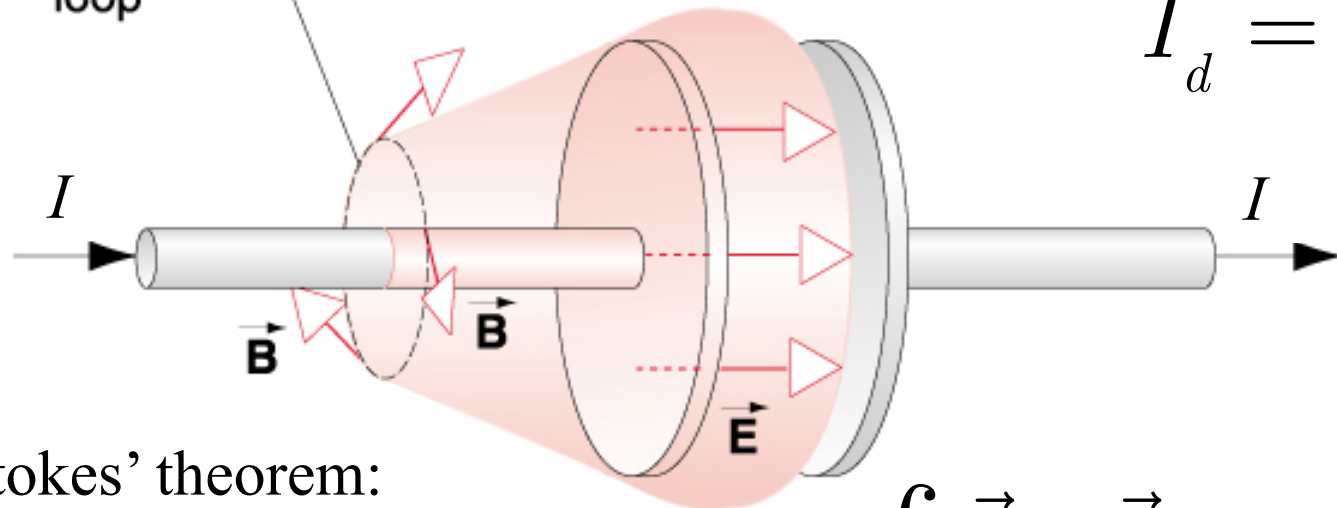
# Maxwell's displacement current



$$\oint \vec{B} \cdot d\vec{s} = \mu_0 (I + I_d)$$

Ampèrian loop

$$I_d = \epsilon_0 \frac{d\Phi_E}{dt}$$



$$\oint \vec{B} \cdot d\vec{l} = \mu_0 I + \mu_0 \epsilon_0 \frac{d\Phi_E}{dt}$$

Stokes' theorem:  
The choice of surface  
should not matter

# Maxwell's equations

**Table 29.2** Maxwell's Equations

Law	Mathematical Statement	What It Says
Gauss for $\vec{E}$	$\oint \vec{E} \cdot d\vec{A} = \frac{q}{\epsilon_0}$	How charges produce electric field; field lines begin and end on charges.
Gauss for $\vec{B}$	$\oint \vec{B} \cdot d\vec{A} = 0$	No magnetic charge; magnetic field lines don't begin or end.
Faraday	$\oint \vec{E} \cdot d\vec{r} = -\frac{d\Phi_B}{dt}$	Changing magnetic flux produces electric field.
Ampère	$\oint \vec{B} \cdot d\vec{r} = \mu_0 I + \mu_0 \epsilon_0 \frac{d\Phi_E}{dt}$	Electric current and changing electric flux produce magnetic field.

The main thing to note here is the symmetry in the last two equations: a time varying magnetic field produces an electric field; similarly, a time varying electric field produces a magnetic field.

# Maxwell's equations in vacuum

$$\Phi_E = \oint \vec{E} \cdot d\vec{A} = 0$$

$$\Phi_B = \oint \vec{B} \cdot d\vec{A} = 0$$

$$\oint \vec{B} \cdot d\vec{l} = \mu_0 \epsilon_0 \frac{d\Phi_E}{dt}$$

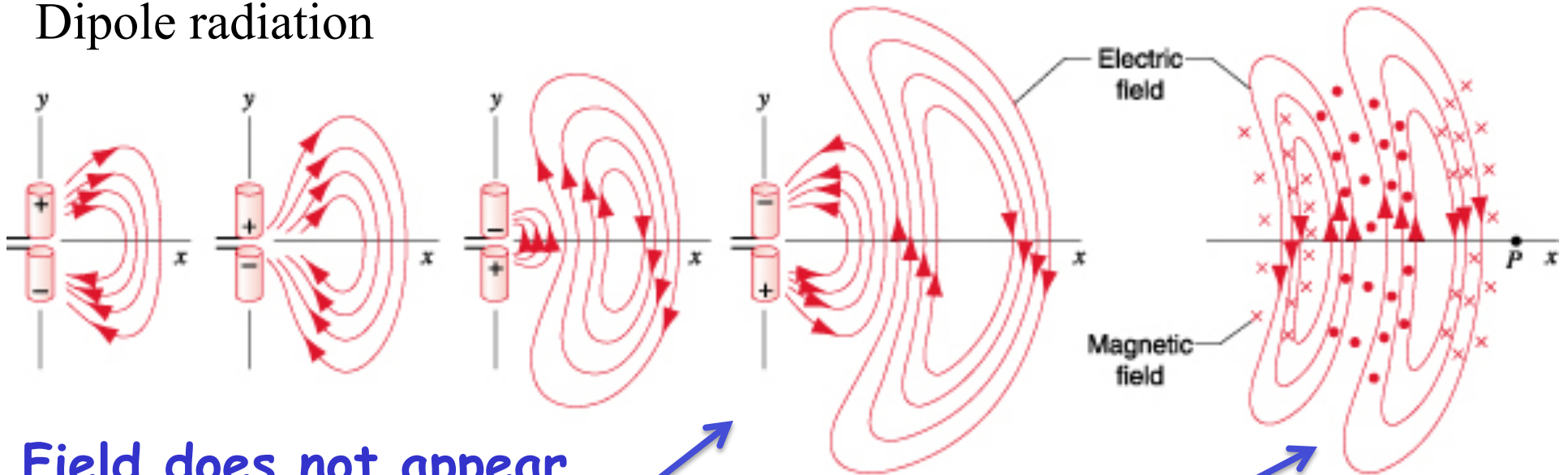
$$\oint \vec{E} \cdot d\vec{l} = -\frac{d\Phi_B}{dt}$$

The main thing to note here is the symmetry in the last two equations: a time varying magnetic field produces an electric field; similarly, a time varying electric field produces a magnetic field.



# Electromagnetic waves

Dipole radiation

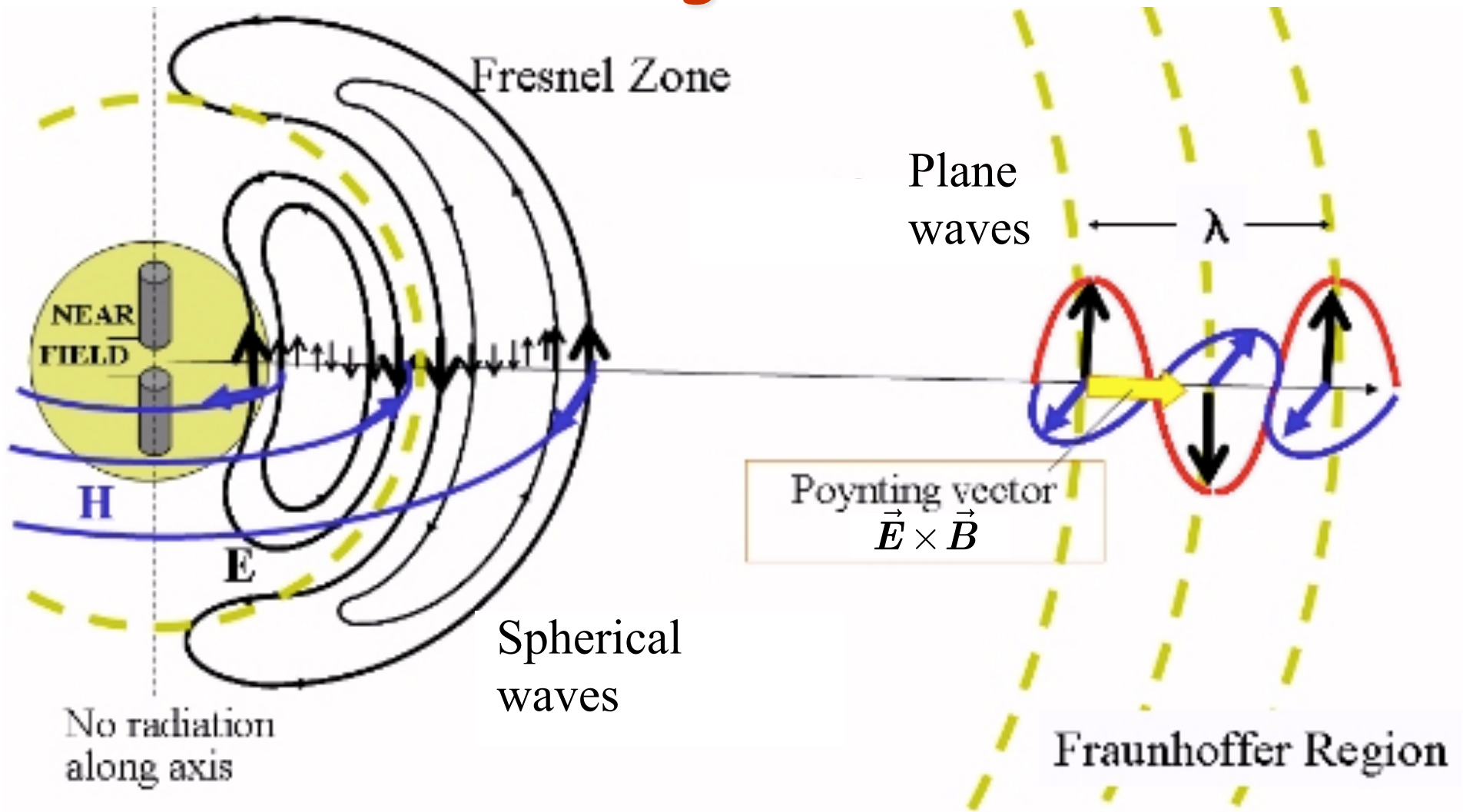


Field does not appear instantaneously

Propagation

Electromagnetic perturbation breaks completely free from the charge/current

# Electromagnetic waves



Maxwell's equations guarantee that electric and magnetic fields are perpendicular to each other and perpendicular to the direction of propagation; **they are polarized.**

# Maxwell's equations in vacuum

$$\Phi_E = \oint \vec{E} \cdot d\vec{A} = 0$$

$$\Phi_B = \oint \vec{B} \cdot d\vec{A} = 0$$

$$\oint \vec{B} \cdot d\vec{l} = \mu_0 \epsilon_0 \frac{d\Phi_E}{dt}$$

$$\oint \vec{E} \cdot d\vec{l} = -\frac{d\Phi_B}{dt}$$

## Stokes' Theorem:

Gives differential form of Maxwell's equations

$$\left[ \nabla \times \vec{B} = \mu_0 \epsilon_0 \frac{d\vec{E}}{dt} \right]$$

$$\left[ \nabla \times \vec{E} = -\frac{d\vec{B}}{dt} \right]$$

# Let there be light!!

## Stokes' Theorem:

Gives differential form  
of Maxwell's equ'ns

$$\nabla \times \vec{B} = \mu_0 \varepsilon_0 \frac{d\vec{E}}{dt}$$

$$\nabla \times \vec{E} = -\frac{d\vec{B}}{dt}$$

Maxwell's equations in vacuum can be solved simultaneously  
to give identical differential equations for  $E$  and  $B$ :

$$\nabla^2 \vec{E} = \mu_0 \varepsilon_0 \frac{\partial^2 \vec{E}}{\partial t^2} \quad \text{and} \quad \nabla^2 \vec{B} = \mu_0 \varepsilon_0 \frac{\partial^2 \vec{B}}{\partial t^2}$$

$$\nabla^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}$$

*The Laplacian differential  
operator*

# The Electromagnetic Wave Equations

# Let there be light!!

## Stokes' Theorem:

Gives differential form  
of Maxwell's equ'ns

$$\nabla \times \vec{B} = \mu_0 \epsilon_0 \frac{d\vec{E}}{dt}$$

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Maxwell's equations in vacuum can be solved simultaneously  
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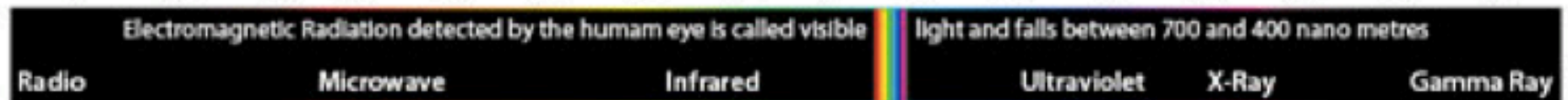
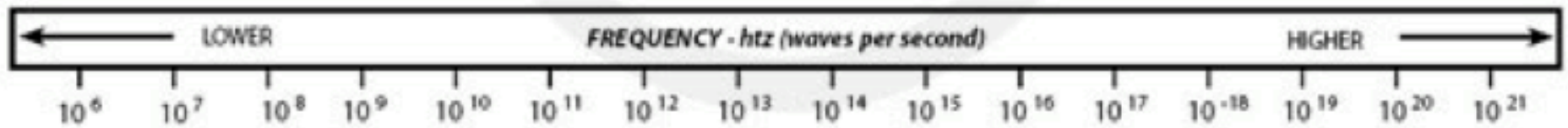
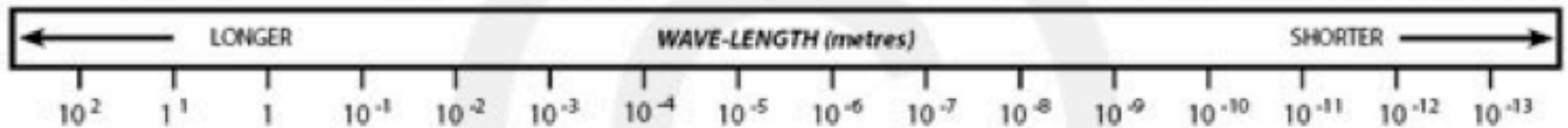
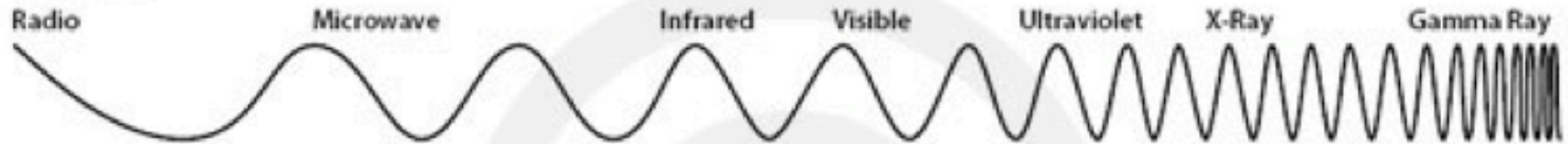
Waves in one-dimension (traveling along  $x$ )

# The Electromagnetic Wave Equations

# THE ELECTRO MAGNETIC SPECTRUM

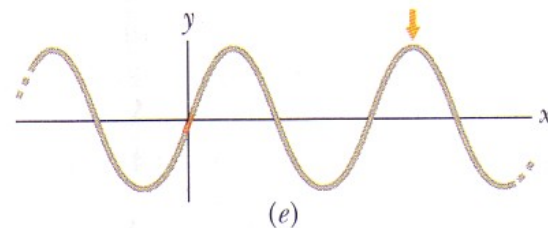
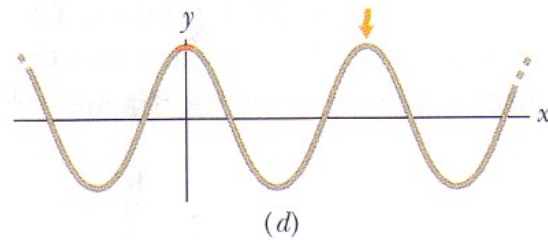
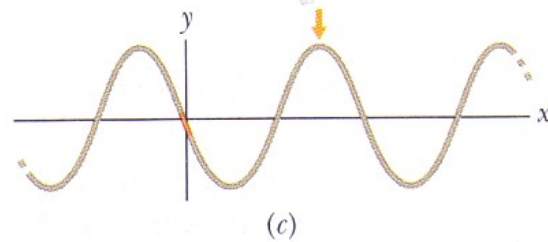
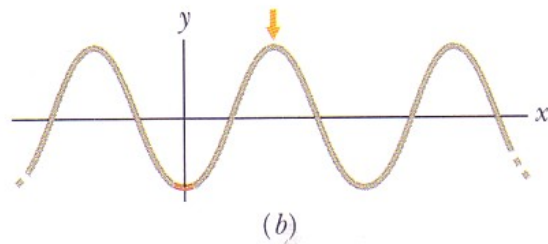
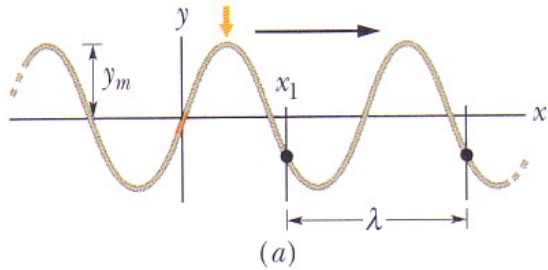
1 metre = 100cm 1 cm = 10mm 1 millimetre = 1000 microns 1 micron = 1000 nanometres (nm) - one nanometer is one billionth of a metre  
 $10^{-5} = 0.00001$   $10^5 = 100,000$

WAVE (type)



# Review of waves (PHY2048)

Transverse wave



Displacement  $y(x,t) = A \sin(kx \pm \omega t + \phi)$

Amplitude  $A$

angular wavenumber  $k$

angular frequency  $\omega$

Phase  $\phi$

Phase shift

$k = \frac{2\pi}{\lambda}$   $k$  is the angular wavenumber.

$\omega = \frac{2\pi}{T}$   $\omega$  is the angular frequency.

frequency  $f = \frac{1}{T} = \frac{\omega}{2\pi}$

velocity  $v = \mp \frac{\omega}{k} = \mp \frac{\lambda}{T} = \mp f \lambda$

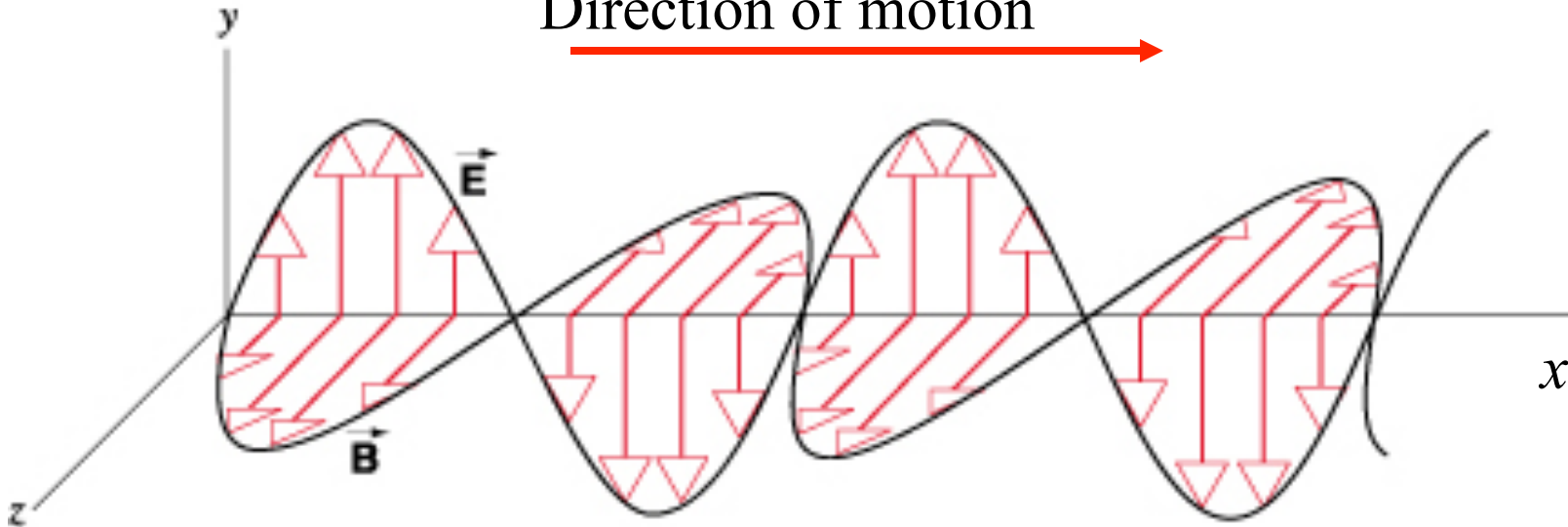
# Electromagnetic waves

- The  $E$  and  $B$  fields are still related via Ampère's and Faraday's laws.
- For a plane wave traveling in the  $x$  direction (see text):

$$\vec{E}(x, t) = E_p \sin(kx - \omega t) \hat{j}$$

$$\vec{B}(x, t) = B_p \sin(kx - \omega t) \hat{k}$$

Direction of motion

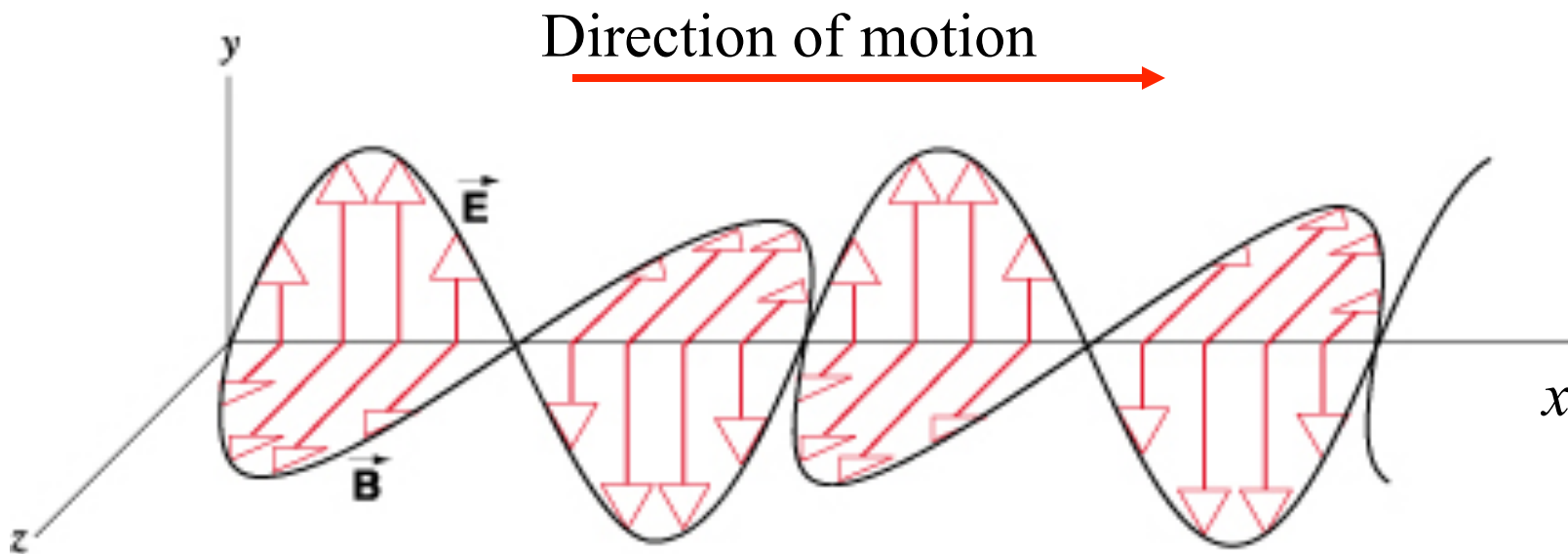




# Electromagnetic waves

- The  $E$  and  $B$  fields are still related via Ampère's and Faraday's laws.
- For a plane wave traveling in the  $x$  direction (see text):

$$\frac{\partial E_z}{\partial x} = \frac{\partial B_y}{\partial t}, \quad \frac{\partial E_y}{\partial x} = -\frac{\partial B_z}{\partial t}, \quad E_x = B_x = 0$$



# Electromagnetic waves

• Plugging these wave solutions into the wave equation:

$$\nabla^2 E_y = -k^2 E_y = \mu_0 \epsilon_0 \frac{\partial^2 E_y}{\partial t^2} = -\omega^2 \mu_0 \epsilon_0 E_y$$

$$\Rightarrow \frac{\omega^2}{k^2} = c^2 = \frac{1}{\mu_0 \epsilon_0}, \quad \text{or} \quad c = \sqrt{\frac{1}{\mu_0 \epsilon_0}}$$

• Plugging these wave solutions into Faraday's law:

$$\frac{\partial E_y}{\partial x} = k E_p \cos(kx - \omega t) = -\frac{\partial B_z}{\partial t} = \omega B_p \cos(kx - \omega t)$$

$$\Rightarrow \frac{E_p}{B_p} = \frac{\omega}{k} = c$$